

# Mueller matrix approach to the propagation of light in stratified anisotropic media

F. Ferrieu<sup>\*1</sup>, E. Garcia-Caurel<sup>2</sup>, M. Stchakovsky<sup>3</sup>

STMICROELECTRONICS, France, <sup>2</sup> LPICM-Ecole polytechnique, France, <sup>3</sup> Horiba Jobin Yvon, France  
[frederic.ferrieu@cea.fr](mailto:frederic.ferrieu@cea.fr), [enric.garcia-caurel@polytechnique.edu](mailto:enric.garcia-caurel@polytechnique.edu)



• Propagation of light in anisotropic material suscites a growing interest of new synthesized materials for nanotechnologies. For a large number of anisotropic substrates or thin films, weak birefringence can be measured by new polarimeters.

• The experimentalist has then to deal with a large number of acquired data such as with these 4x4 Mueller matrix.

• The aim is to revise the existing formalism detailed in several publications and to give an overall view in order to build a general computational algorithm to calculate the Mueller matrix and the Jones matrix of a given sample.

• The explicit equations given hereafter can be applied for both transmitted and reflected light propagating into anisotropic multilayered media. They can be implemented with the aid of commercial linear algebra libraries.

## Evaluation of $\exp(ik_0\Delta d_i)$

The exponential matrix applying the Cayley Hamilton theorem.

$$T_p = \exp ik_0 \Delta d = \beta_0 I + \beta_1 \Delta + \beta_2 \Delta^2 + \beta_3 \Delta^3 \quad \exp\{ik_0 q_k d\} \equiv \sum_{j=0}^3 \beta_j q_k^j \dots k=1, \dots, 4$$

$$\beta_0 = \sum_{i=1}^4 q_j q_k q_i \frac{f_i}{q_{ij} q_{ik} q_{ii}} \dots \quad \beta_2 = -\sum_{i=1}^4 (q_j + q_k + q_i) \frac{f_i}{q_{ij} q_{ik} q_{ii}} \dots$$

$$\beta_1 = \sum_{i=1}^4 (q_j q_k + q_j q_i + q_k q_i) \frac{f_i}{q_{ij} q_{ik} q_{ii}} \quad \beta_3 = \sum_{i=1}^4 \frac{f_i}{q_{ij} q_{ik} q_{ii}} \dots$$

$e^{iq_k k_0 d} = M\{\beta_0, \beta_1, \beta_2, \beta_3\}^T$  The exponential of the  $\Delta$  matrix is a function of the eigenvalues of the same  $\Delta$  matrix.

The value of the eigenvalues  $\beta_i$  can be also found using commercially available software to solve the corresponding system of linear equations.

## Propagation of light: Maxwell Equations

The propagation of light through a general linear and non-magnetic anisotropic medium can be described using Maxwell equations written in their wave form:

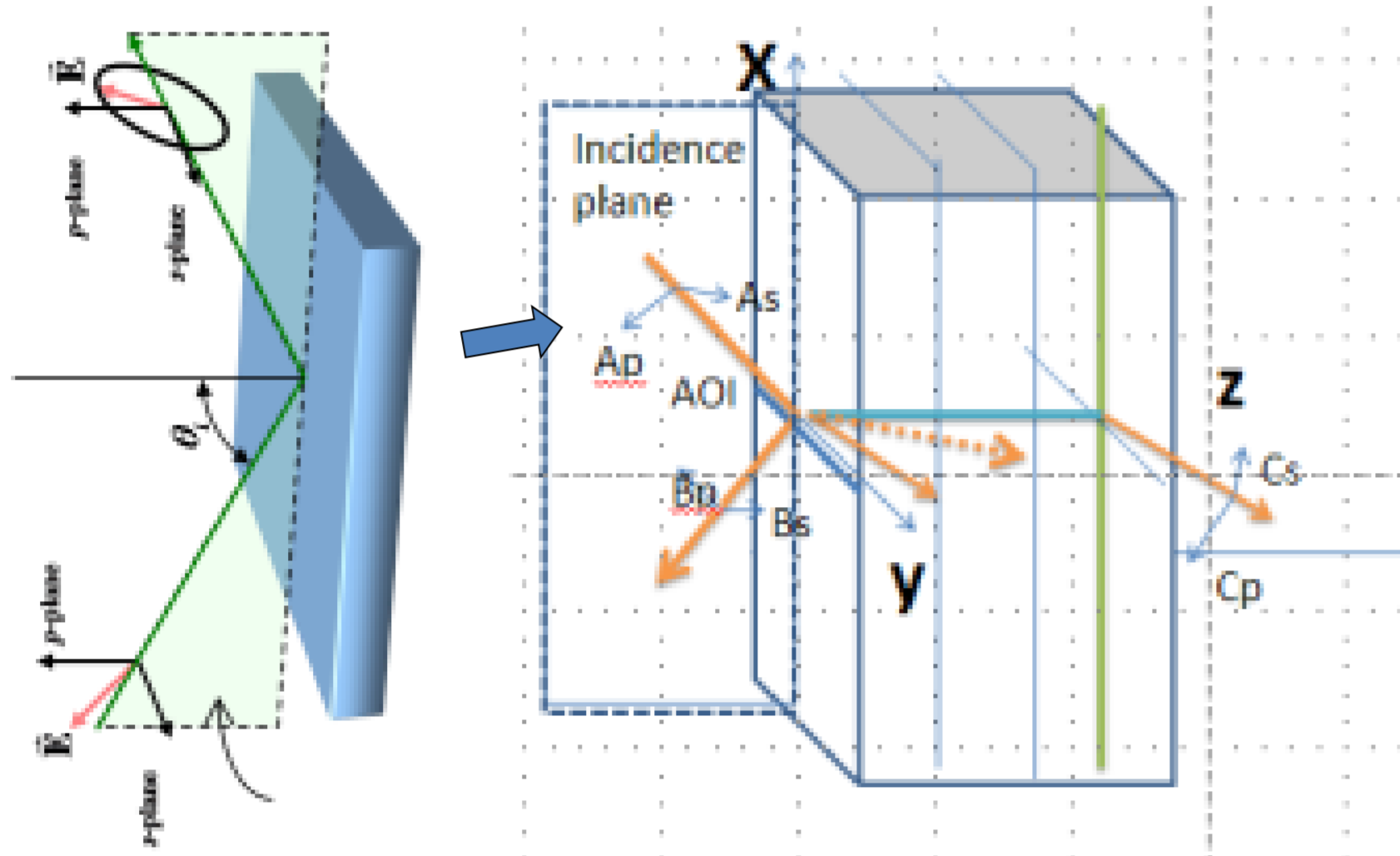
$$\text{curl } \vec{H} = j\omega \vec{D}$$

$$-\text{curl } \vec{E} = j\omega \vec{B}$$

$$\frac{\partial \Psi(z)}{\partial z} = -j\omega \Delta \Psi(z)$$

$$\Psi = [E_x, E_y, H_x, H_y]^T$$

$$k_0 = \omega/c$$



$$\Psi(z + d_i) = \exp\{ik_0 \Delta d_i\} \Psi(z) = T_{pi} * \Psi(z)$$

$$\begin{bmatrix} A_s \\ B_s \\ A_p \\ B_p \end{bmatrix} = T \begin{bmatrix} C_s \\ D_s \\ C_p \\ D_p \end{bmatrix} = \begin{bmatrix} T_{11} & T_{12} & T_{13} & T_{14} \\ T_{21} & T_{22} & T_{23} & T_{24} \\ T_{31} & T_{32} & T_{33} & T_{34} \\ T_{41} & T_{42} & T_{43} & T_{44} \end{bmatrix} \begin{bmatrix} C_s \\ 0 \\ C_p \\ 0 \end{bmatrix}$$

Multilayer case ( thickness  $d_i$ )

$$T = L^{-1}_a \prod_{i=1}^n T_{pi}(-d_i) L_f$$

$$L_a, L_f$$

Transition matrices respectively between air and film and film with substrate

## Evaluation of Transition matrices air/film $L_a, L_f$

The continuity of tangential components imposes the relation between, **E**, and **H**.

The projection of the fields on both, s and p components yield:

$$L_a \Psi_a = \Psi_{inc}(z=0) + \Psi_{ref}(z=0) \quad \left\{ \begin{array}{l} \Psi_{inc} = [A_p \cos \Phi_a, A_s, -n_a A_s \cos \Phi_a, n_a A_p] \\ \Psi_{ref} = [-B_p \cos \Phi_a, B_s, n_a B_s \cos \Phi_a, n_a B_p] \end{array} \right.$$

$$\{L_{ij}\}^{-1} = \begin{bmatrix} 0 & 1 & -1/n_a \cos \Phi_a & 0 \\ 0 & 1 & 1/n_a \cos \Phi_a & 0 \\ 1 & 0 & 0 & 1/n_a \\ -1 & 0 & 0 & 1/n_a \end{bmatrix}$$

$\Phi_a$  AOI (Angle of Incidence)  
 $n_a$ =ambient index  $n_a=1$

## Transition matrices film/anisotropic substrate $L_f$

$\Psi_f$  is a linear combination of the eigenvectors  $\Xi$  of  $\Delta$

$$\Psi_f = \sum_{i=1}^4 c_i \Xi_i(q_i) \quad \{L_{fjk}\} = \Xi_j, L_{fj(k+1)} = 0 \quad j=1, \dots, 4 \quad k=1, 3 \quad \text{Re}\{q_i\} > 0$$

For a biaxial system eigenvectors are easily obtained

Such as with axis aligned. To the laboratory frame one gets thus for the  $c_i$ :

$$\{c_i\} = \{0, -n_y \cos \Phi_y, 0, n_x\} \quad (L_{fjk}) = \begin{bmatrix} \Xi_{11} = q_{21} & 0 & \Xi_{13} = q_{31} & 0 \\ \Xi_{21} = q_{22} & 0 & \Xi_{23} = q_{32} & 0 \\ \Xi_{31} = q_{23} & 0 & \Xi_{33} = q_{33} & 0 \\ \Xi_{41} = q_{24} & 0 & \Xi_{43} = q_{34} & 0 \end{bmatrix}$$

## Euler angles

In general the coordinate system of the laboratory and the main coordinate system of the anisotropic media are not the same.

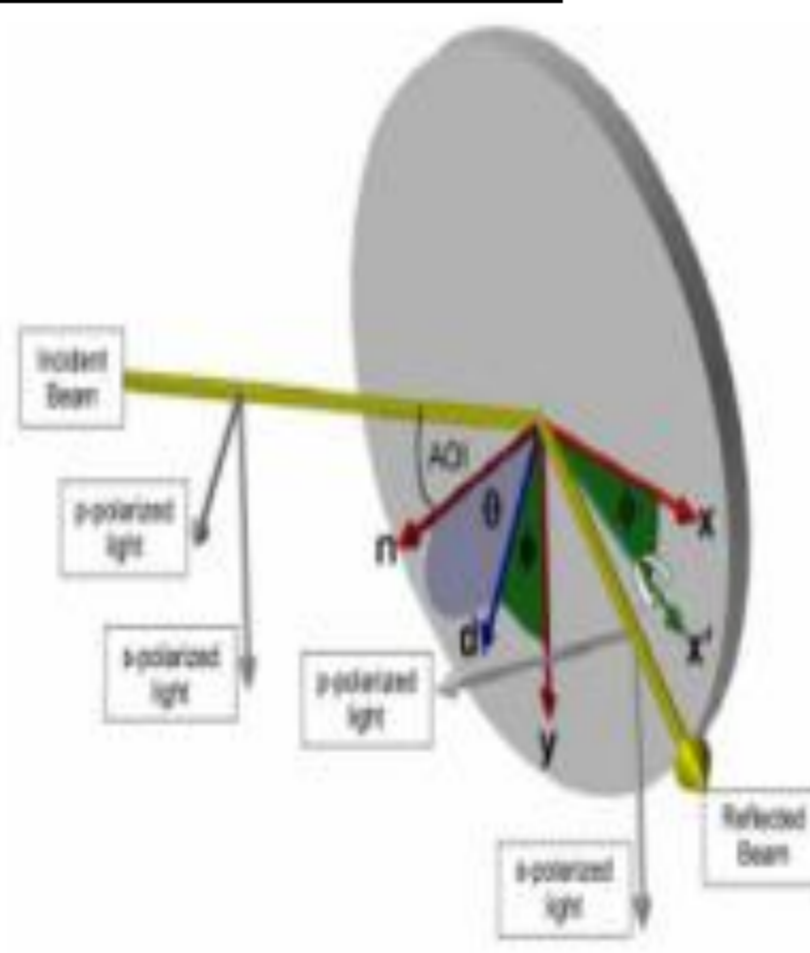
They can be related by a rotation transformation. A common way of writing a rotation transformation is to use the Euler matrices and the associated Euler angles  $(\theta, \Psi, \Phi)$ .

$$\varepsilon = \begin{bmatrix} \varepsilon_x = \varepsilon_{11} & 0 & 0 \\ 0 & \varepsilon_y = \varepsilon_{22} & 0 \\ 0 & 0 & \varepsilon_z = \varepsilon_{33} \end{bmatrix} \rightarrow \{\varepsilon_{ij}\} = \begin{bmatrix} \varepsilon_{11} & \varepsilon_{12} & \varepsilon_{13} \\ \varepsilon_{21} & \varepsilon_{22} & \varepsilon_{23} \\ \varepsilon_{31} & \varepsilon_{32} & \varepsilon_{33} \end{bmatrix}$$

Main frame of the medium

Laboratory frame

$$\{\varepsilon_{ij}\} = A(\varphi, \psi, \theta) \varepsilon A^{-1}(\varphi, \psi, \theta)$$



## Calculation of Mueller matrices

The  $T_{ij}$  elements correspond to the general transfer matrix= matrix product of each layer thickness  $d_i$

The four elements of the Jones matrix are calculated through the Kronecker product with matrices  $A$  and  $A^{-1}$  (P. Yeh)

$$J \equiv \begin{bmatrix} r_{pp} & r_{ps} \\ r_{sp} & r_{ss} \end{bmatrix} \quad M \equiv A (J \otimes J^*) A^{-1} \quad A \equiv \begin{bmatrix} 1 & 0 & 0 & 1 \\ 1 & 0 & 0 & -1 \\ 0 & 1 & 1 & 0 \\ 0 & -j & j & 0 \end{bmatrix}$$

$$r_{ps} = \left( \frac{B_s}{A_p} \right)_{A_{z2}=0} \equiv \frac{T_{11}T_{23} - T_{21}T_{13}}{T_{11}T_{33} - T_{13}T_{31}} \quad r_{ss} = \left( \frac{B_s}{A_s} \right)_{A_p=0} \equiv \frac{T_{21}T_{33} - T_{23}T_{31}}{T_{11}T_{33} - T_{13}T_{31}}$$

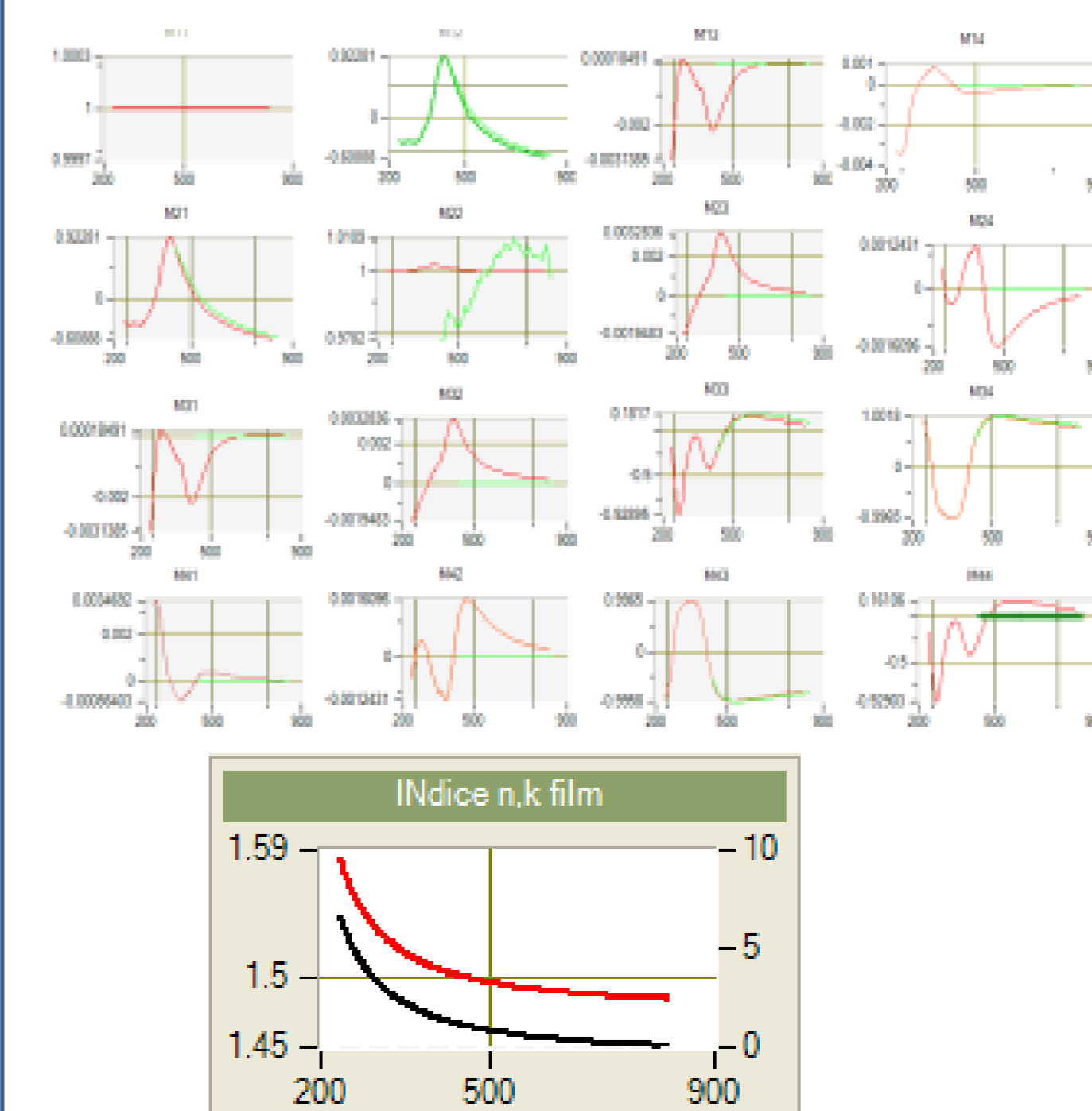
$$r_{pp} = \left( \frac{B_p}{A_p} \right)_{A_{z2}=0} \equiv \frac{T_{11}T_{43} - T_{41}T_{13}}{T_{11}T_{33} - T_{13}T_{31}} \quad r_{sp} = \left( \frac{B_p}{A_s} \right)_{A_p=0} \equiv \frac{T_{41}T_{33} - T_{43}T_{31}}{T_{11}T_{33} - T_{13}T_{31}}$$

## Construction of the $\Delta$ Matrix :

Following the curl operator in Maxwell equations for **E** and **H** the  $\Delta$  matrix depends on the dielectric tensor and the component  $k_x$  of the wave propagation vector  $k_0$ ,

$$\Delta = \begin{bmatrix} -k_x \frac{\varepsilon_{31}}{\varepsilon_{33}} & -k_x \frac{\varepsilon_{32}}{\varepsilon_{33}} & 0 & 1 - \frac{k_x^2}{k_0^2} \\ 0 & 0 & -1 & 0 \\ \varepsilon_{23} - \varepsilon_{13} \frac{\varepsilon_{31}}{\varepsilon_{33}} & k_x^2 - \varepsilon_{22} + \varepsilon_{23} \frac{\varepsilon_{32}}{\varepsilon_{33}} & 0 & k_x \frac{\varepsilon_{23}}{\varepsilon_{33}} \\ \varepsilon_{11} - \varepsilon_{13} \frac{\varepsilon_{31}}{\varepsilon_{33}} & \varepsilon_{12} - \varepsilon_{13} \frac{\varepsilon_{32}}{\varepsilon_{33}} & 0 & -k_x \frac{\varepsilon_{13}}{\varepsilon_{33}} \end{bmatrix} \quad k_x = n_a \sin \Phi_a$$

## Simulation software



Here are some snapshots showing several interfaces of our software: Input parameters, Dispersion of index, Mueller matrix.

